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SERGEI GUKOV

SHING-TUNG YAU

ERIC ZASLOW

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Duality and Fibrations on G_2 Manifolds

Sergei Gukov, Shing-Tung Yau, and Eric Zaslow

Abstract

We argue that G_2 manifolds for M-theory admitting string theory Calabi-Yau duals are fibered by coassociative submanifolds. Dual theories are constructed using the moduli space of M-five-brane fibers as target space. Mirror symmetry and various string and M-theory dualities involving G_2 manifolds may be incorporated into this framework. To give some examples, we construct two non-compact manifolds with G_2 structures: one with a $K3$ fibration, and one with a torus fibration and a metric of G_2 holonomy. Kaluza-Klein reduction of the latter solution gives abelian BPS monopoles in $3 + 1$ dimensions.

1. Introduction

One of the main achievements in string theory during the last decade was the discovery of string dualities and relations among them. A particularly rich and interesting example of string duality is *mirror symmetry* between pairs of Calabi-Yau manifolds. A geometric framework for understanding this duality was proposed in [32], and involves constructing the mirror manifold by dualizing a torus fibration. This construction arose from the correspondence among nonperturbative states of dual theories. M-theory has united the disparate string theories and promises to reveal the nature of string dualities. In M-theory, the analogue of a Calabi-Yau manifold is a manifold with G_2 holonomy, simply by the counting of dimensions: what was $10 = 4 + 6$ for string theory is $11 = 4 + 7$ in M-theory. According to this simple formula, seven-dimensional G_2 -holonomy manifolds are natural candidates for minimally supersymmetric (and phenomenologically interesting [3]) compactifications of M-theory to $3 + 1$ dimensions. If manifolds with G_2 holonomy are M-theory analogues of Calabi-Yau spaces, then what is the corresponding notion of mirror symmetry, and what is the geometry behind duality? Is there a fibration structure on G_2 manifolds relevant to this and possibly other string/M-theory dualities? These are the questions that one might naturally ask, and that we attempt to address in this paper.

We argue that, just as Calabi-Yau manifolds involved in mirror symmetry are fibered by special Lagrangian three-tori, in M-theory G_2 -holonomy manifolds which admit string theory duals are fibered by coassociative four-manifolds. M-Theory on a seven-manifold X , with G_2 holonomy, leads to an effective field theory in four dimensions with $\mathcal{N} = 1$ supersymmetry. The same is true for the heterotic string theory on a Calabi-Yau manifold,

OV Z LO

Y,

¹

Z,

fl

fl

² X, Y X, Z
X

Y Z,

G_2
 G_2

C.

C

Y Z,

C.

K

T^3 ,

K

K

G_2

T^7

Remark 1.

O \xrightarrow{j} G_2 \xrightarrow{f} b \xrightarrow{p} p
 \cdot F \xrightarrow{x} p \xrightarrow{f} C \xrightarrow{b} $f G_2$ \xrightarrow{f} per se
 K \xrightarrow{z} K \xrightarrow{fib} f \xrightarrow{ff} G_2 \xrightarrow{f} G_2 \xrightarrow{f} C \xrightarrow{b} f
 \cdot O \xrightarrow{f} xp \xrightarrow{f} G_2 \xrightarrow{f} A
 \cdot f $\xrightarrow{f G_2}$ Vf \cdot f \xrightarrow{fib} f \cdot
 \cdot G_2 \xrightarrow{f} p \xrightarrow{fib} F \xrightarrow{M} f
 \cdot p $\xrightarrow{f G_2}$

Remark 2. R

\cdot W \xrightarrow{p} \cdot 2 \xrightarrow{f} \cdot 3 $\xrightarrow{4}$ b \xrightarrow{p} p
 x b \cdot p \cdot b \xrightarrow{p} \cdot \cdot 5

¹One requires $E \rightarrow Y$ to obey $p_1(E) = p_1(TY)$ and $c_1(E) = 0$ so that there is no anomaly, *i.e.* the heterotic theory contains no fivebranes.

²Kachru and Vafa first found heterotic-Type-IIA (Y, Z) pairs in [19]; some (X, Z) pairs are studied in [4, 5].

2. Fibrations from Brane Moduli Spaces

2.1. Fibrations from M-theory



³We recall that the field content of M-theory contains a three-form H with four-form field strength; it obeys $dH = \delta_D$, where D is the five-brane world-volume. This leads to a condition that the normal bundle have trivial Euler characteristic, which is true for the examples $(T^4, K3)$ in this paper.

⁴One could also consider the Kaluza-Klein modes of the two-form, say, with only one leg along Σ . These modes, however, give rise to $U(1)$ gauge vector fields on Σ and, therefore, are non-dynamical.

$$b_2^+ C$$

$$C b_2^+ C$$

$$H^2 C \mathbb{R} / H^2 C \mathbb{Z} . \quad p$$

D

$$H^3 D \mathbb{R} / H^3 D \mathbb{Z}$$

$$X \quad b_2^+ C \quad b_2^- C \quad C \quad T^4 \quad C \quad K . \quad b_2^+ C \quad H^2 C, \mathbb{Z} .$$

$$C \sim T^4 \quad C \quad b_2^+ \quad b_2^- \quad , \quad b_2^+ C \quad C \quad \text{fl}$$

$$C \sim K \quad b_2^+ \quad , \quad b_2^- \quad ,$$

$$b_2^+ C \quad C \quad T^4 \quad K ,$$

C B

C .

⁵The metric is typically derived from the second-order variation of the five-brane action with respect to world-sheet time-dependence of the moduli fields, and the fermionic terms are derived from supersymmetry.

OV Z LO

$C \sim T^4,$

G_2

$H^2 C \mathbb{R}$

G_2 C
 fl fl

$C \sim T^4 \quad C \sim K$

2.2. Torus Fibrations and Type-II/M-theory Duality

T^4

$C \sim T^4$

X

6

different

$X,$

mirror

dual

G_2

v

$\mathbb{R}^4/\mathbb{R}v$

G_2 $\mathbb{R}^4,$ G_2 \mathbb{R}^4
 0 \mathbb{R}^7

${}^2_+ \mathbb{R}^4 \sim \mathbb{R}^4 \perp,$

⁶If the fibration of T^4 over T^3 is not changing, then the dilaton is constant and $X \cong CY \times S^1$, so we recover $\mathcal{N} = 2$ supersymmetry.

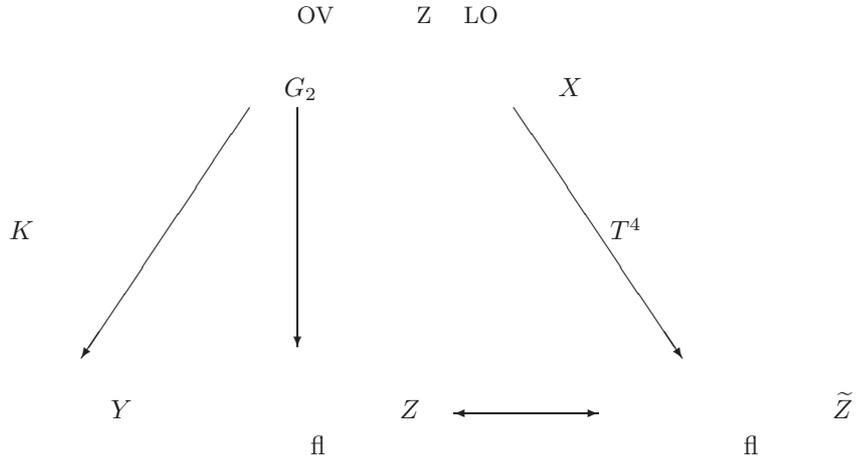


Figure 1. V

$$v, \quad w \in \mathbb{R}^4 / \mathbb{R}v, \quad n$$

$$0 v, w, n,$$

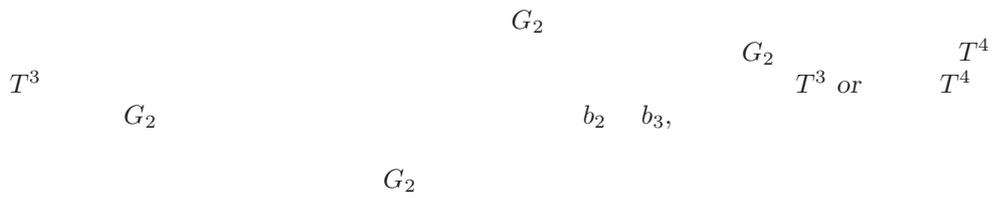
7

G_2 X ,
brane reduction is mirror to Kaluza-Klein reduction.

C

B

..

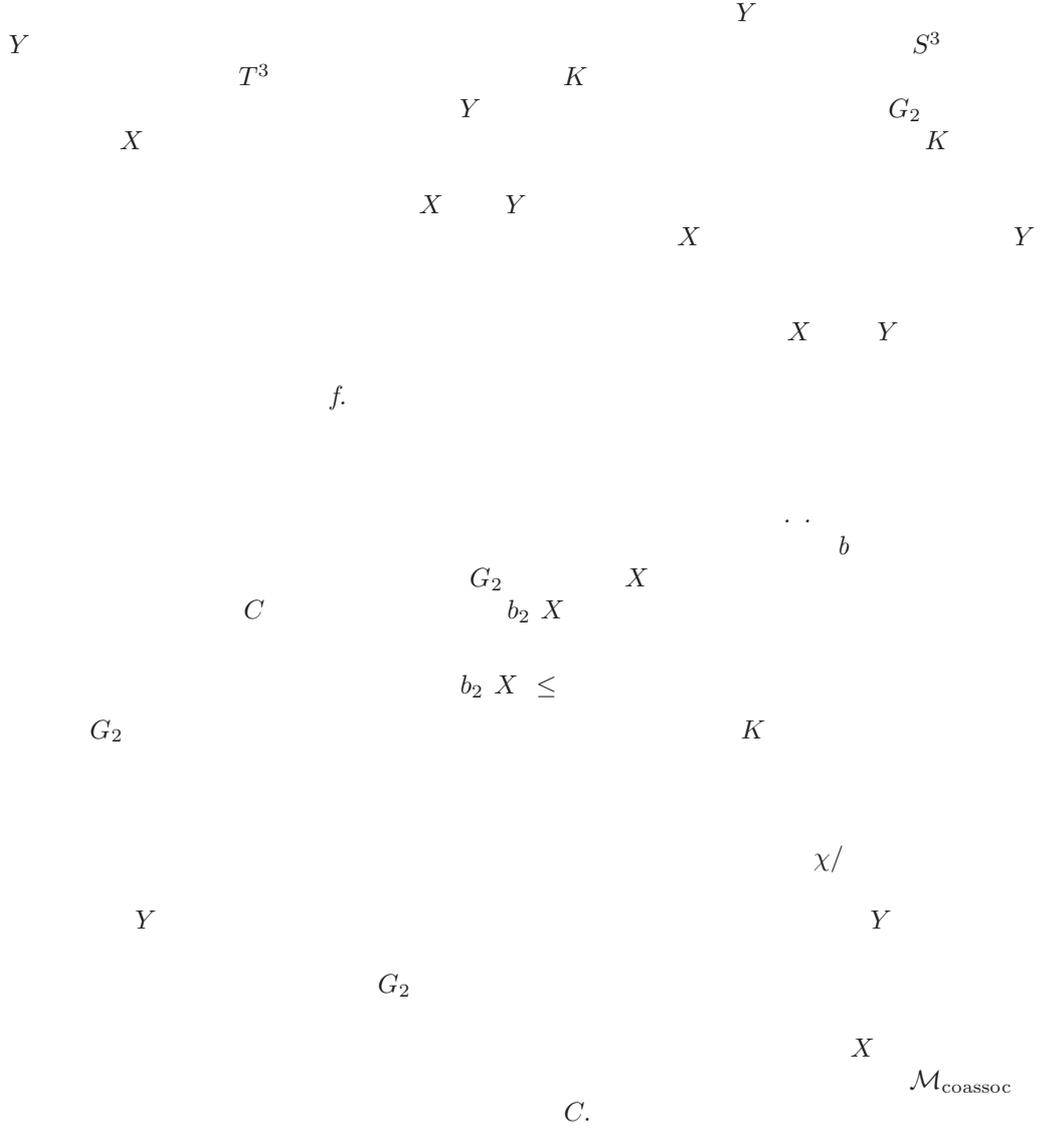


⁷For example, if we write

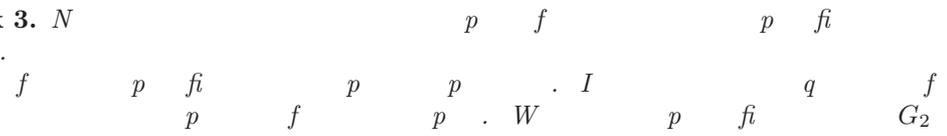
$$\Phi_0 = e_{125} + e_{345} + e_{136} - e_{246} + e_{147} + e_{237} + e_{567} \tag{2.2}$$

and $v = e_1$, then the duality pairs $e_2 \leftrightarrow e_5$, as $\iota_v \Phi_0 = e_{25} + e_{36} + e_{47}$, etc.

2.3. K Fibrations and Heterotic/M-theory Duality



Remark 3. N



K
 X G_2 X X
 K p, q q X B p
 \mathcal{E} \mathcal{E}_p X K q $v \mathcal{E}_p$
 $v \mathcal{E}_p$ $ch \mathcal{E}_p \sqrt{A K}$
 v $, , -$
 X K K $/$
 $q \in B$ K X $b_2^+ K$ $,$
 X
 v_{dual} $, , \cdot$
 v $, , -$ K $p_1 K$
 v $, ,$
 \mathcal{E} v r, l, s
 $dim M \mathcal{E} = l^2 - rs$ $.$
 v $, , \cdot$
 K T^4
 T^4 T^4 $b_1 T^4$ $b_3^- T^4$
 fl

$$\begin{array}{ccccccc}
 R & & f & B & ? & & \\
 & & K & & T^4 & & \\
 X & & & & B & & X \\
 & & B & & & & \\
 & & \text{fl } U & & X & G_2 & b_1 X
 \end{array}$$

3. A K Fibration

3.1. Idea and Basic Set-Up

$$\begin{array}{ccccccc}
 & & G_2 & & K & & S^3, \\
 & & , & & & & \\
 & & G_2 & & K & & S^3 \setminus X \\
 & & & & g_B, & & B \text{ } SO, /SO \\
 X & B \times K & & & & & SO, . \\
 & & & & \pi \text{ } X \rightarrow B & & \\
 & & & & & & p \in X \\
 & & & & & & T_V X \text{ } TX, \\
 & & & & \text{connection,} & & \\
 T_H X & TX. & & & TX & TX & T_H X \oplus T_V X, \\
 P_H & P_V & & & & & \\
 & & & & b \in B & & K \\
 & & & & & & K \\
 & & & & & & \mathcal{M}_{K3} \text{ } SO, /SO \times SO \\
 X & \dots & & & SO, \mathbb{Z} & & B \text{ } \mathcal{M}_{K3}. \\
 & & & & K & &
 \end{array}$$

⁹This construction works for T^4 fibrations too, if in the following we simply replace $SO(3, 19)$ and its maximal compact subgroup by $SO(3, 3)$ with its corresponding subgroup.

$$\tau: B \rightarrow \mathcal{M}_{K^3}, \quad \text{fl} \quad \tau, \quad \text{fl}$$

3.2. The Fiberwise Metric and Connection

$$\begin{array}{c} RicMet \\ K \end{array} \xrightarrow{p} \begin{array}{c} RicMet \times K \\ q \\ K \end{array} \xrightarrow{Diff} \begin{array}{c} p, q \\ RicMet \\ RicMet \times_{Diff} K \\ \text{fl } K \end{array}$$

$$\begin{array}{c} \mathcal{M}_{K^3} \\ \text{fl } K \end{array} \xrightarrow{RicMet/Diff} \begin{array}{c} X \\ \text{fl } K \end{array} \xrightarrow{\tau} \begin{array}{c} X \rightarrow B \\ \pi^{-1} b \end{array}$$

$$b \in B, \quad U \subset B, \quad TX, \quad V_b \subset X, \quad \pi^{-1} U \sim U \times K, \quad \gamma: \pi^{-1} b \rightarrow V_b$$

$$\text{fixed } g_t, \quad f_t: K \rightarrow K, \quad \gamma_t, \quad \frac{d}{dt} f_t^* g_t \perp$$

$$\Gamma_t: V \rightarrow V, \quad \Gamma_t(\gamma_t, f_t q), \quad V_H \equiv \frac{d}{dt} \Gamma_t|_{t=0}$$

Lemma 5. (3.11) $q \rightarrow f_t$

$$P \xrightarrow{f} \xi, \quad \xi = \xi_\mu dx^\mu, \quad \eta = \frac{d}{dt} g_t|_{t=0}$$

$$A_{\mu\nu} \equiv \frac{d}{dt} f_t^* g_{t \mu\nu}|_{t=0} = \eta_{\mu\nu} - \nabla_\mu \xi_\nu - \nabla_\nu \xi_\mu, \quad \nabla^\mu A_{\mu\nu} = 0$$

$$\int \theta_V^2 / |V|^2. \tag{12}$$

$$\theta : T_b B \rightarrow H_+^2 \pi^{-1} b ,$$

$$\mathfrak{h}_g \sim so \stackrel{p_1 \circ \tilde{\tau}}{\sim} so H_+^2 K \sim H_+^2 K .$$

$$\theta \quad H_+^2$$

$$T_{[g]}^* B \stackrel{p \in X}{\sim} \mathfrak{g}/\mathfrak{h}_g, \quad \pi p \quad g \in B. \quad \{e_i\} \quad \pi^{-1} b \sim K .$$

$$Vol_B \quad \sum_i e_i \wedge \theta_i,$$

$$Vol_B \quad B \quad X$$

$$U, V, W \quad Vol_B \pi_* U, \pi_* V, \pi_* W \quad \theta_{\pi_* U} P_V V, P_V W \quad .$$

$$\tilde{\tau} \quad e_i$$

$$\theta_i$$

3.4. Explicit Formulas

$$\eta_{\mu\nu}$$

Lemma 7. η p .

$$P \quad f. \quad \nabla^\mu \eta_{\mu\alpha} \quad R^\mu_{\alpha\mu\beta}$$

$$\nabla^\mu \nabla_\mu \eta_{\alpha\beta} \quad R_{\alpha}^{\mu} \beta^{\nu} \eta_{\mu\nu} \quad \nabla_\alpha \nabla_\beta \eta^\mu_{\mu} \quad .$$

$$g^{\alpha\beta}$$

$$\nabla^\mu \nabla_\mu \eta^\alpha_{\alpha} \quad .$$

¹²If we write an element C_V of $so(3)$ as $aX + bY + cZ$, then (a, b, c) defines the axis of rotation in three-space. After equating the three-plane \mathbb{R}^3 isometrically with $so(3)$, then θ_V is simply (a, b, c) . That is, $End(E) \cong E$ for oriented, three-dimensional metric vector spaces.

$$\Delta \eta^\alpha_\alpha \int_C C \cdot Vol, \quad \eta^\alpha_\alpha \quad C, \quad \eta$$

$$S^a, a, \dots, S^a \quad \blacksquare \quad i, j, k,$$

$$\eta \quad A^a$$

$$A^a_{\mu\nu} = S^a_{\mu\sigma} \eta^\sigma_\nu - \eta_{\mu\sigma} S^{a\sigma}_\nu.$$

$$A^a$$

$$\eta_{\mu\nu} = \sum_a A^a_{\mu\sigma} S^{a\sigma}_\nu.$$

η 13

$$, \quad H^2_+ K, \quad \tilde{\tau} \quad W \quad \mathbb{R}^{3,19} \sim H^2 K$$

$$S^a, \quad A^a \quad r^a \beta \quad a, \dots, \quad \beta \quad W \quad SO, \quad \sim SO W.$$

$$\eta, \quad A^a, \quad V \quad B$$

$$\theta_V = \sum_a r^a S^a.$$

3.5. d /

$$\tau \quad G_2,$$

$$d \quad \lambda^* \quad G_2 \quad G_2$$

$$G \quad B, \quad \tilde{\tau} \quad SO, \quad \mathbb{Z}$$

$$SO, \quad \mathbb{Z} \quad SO, \quad \mathbb{Z}$$

¹³This depends on some nice facts, including the following identity. Let A be a traceless, symmetric, four-by-four matrix acting on the quaternions \mathbb{R}^4 . Let I, J, K be matrices representing multiplication by i, j, k . Then $A = IAI + JAJ + KAK$.

4. A Torus Fibration

4.1. Outline of Hitchin's Construction

$$\begin{array}{ccc} & G_2 & SU \\ \text{fl} & & G_2 \end{array}$$

$$G_2 \qquad G_2$$

Theorem 8. $f : L \rightarrow M$ $b \in \mathcal{A} \times \mathcal{B}$ $\rho, \sigma \in H^3(M, \mathbb{R})$ $\mathcal{A} \in H^3(M, \mathbb{R})$ $\mathcal{B} \in H^4(M, \mathbb{R})$ b fix
 H fl f f :

$$H = V\rho - V\sigma.$$

$H = V\rho - V\sigma$ b f (f, f, b) ϕ
 $\omega \wedge \rho = \int_M \phi$. If $f : M \rightarrow X$ $t = t_0$, $\rho = \sigma$ f, f, p, b
 $\phi = \rho - \sigma$ (σ, ω^2) f
 $dt \wedge \omega = \rho$,

$$f : G_2 \rightarrow X = M \times a, b \quad a, b.$$

Definition 9. $f : L \rightarrow M$ b f f (n) V TM . $\rho \in {}^p V^*$ b f p b f (n) $GL(V)$ ${}^p V^*$.
 $\rho \in GL(V)$

$$\rho = \omega \quad \omega^{n/2} / M$$

$$\sigma \in {}^4 V^* \sim {}^2 V \otimes {}^6 V^*.$$

$$\sigma^3 \in {}^6 V \otimes {}^6 V^* \sim {}^6 V^* \otimes {}^6 V^*$$

$$V\sigma = \int_M |\sigma^3|^{\frac{1}{2}}.$$

OV Z LO

$$V \rho \quad \rho \in {}^3V^*$$

$$K_\rho : V \rightarrow V \otimes {}^6V^*,$$

$v \in V \quad TM$

$$K v \quad \iota v \rho \wedge \rho \in {}^5V^* \sim V \otimes {}^6V^*.$$

$$\text{tr } K^2 \in {}^6V^{*2}.$$

SL, \mathbb{C}

$$\text{tr } K^2 <$$

$$V \rho \quad \int_M |\sqrt{-\text{tr } K^2}|.$$

$$\mathcal{A} \times \mathcal{B} \sim \Omega_{exact}^3 M \times \Omega_{exact}^4.$$

$$\omega \quad \rho_1, \sigma_1, \rho_2, \sigma_2 \quad \langle \rho_1, \sigma_2 \rangle - \langle \rho_2, \sigma_1 \rangle,$$

$$\rho \quad d\beta \in \Omega_{exact}^p M \quad \sigma \quad d\gamma \in \Omega_{exact}^{n-p+1} M$$

$$\langle \rho, \sigma \rangle = \int_M d\beta \wedge \gamma - \int_M \beta \wedge d\gamma.$$

fl

$$dt \wedge \omega \quad \rho$$

$$d \quad , \quad d^* \quad .$$

G_2

$$v, w \in W, \quad \begin{matrix} W \\ W \end{matrix} \quad TX$$

${}^7W^*$

$$B_\Phi \quad -\iota v \wedge \iota w \wedge .$$

$$K_\Phi : W \rightarrow W^* \otimes \wedge^7 W^*. \quad G_2$$

$$g_\Phi(v, w) = B_\Phi(v, w) \quad K_\Phi^{-\frac{1}{9}}.$$

4.2. Equations for the Metric

$$\begin{aligned}
 & M = S^3 \times T^3. \\
 & \text{SU} \quad \text{SU} \quad M \\
 & \text{SU} \quad G_2 \\
 & \quad \rho \quad \sigma \\
 & \begin{matrix} 1 \\ 2 \\ 3 \end{matrix} \quad \begin{matrix} \psi d\theta & \psi & \theta d\phi, \\ -\psi d\theta & \psi & \theta d\phi, \\ d\psi & \theta d\phi. \end{matrix} \\
 & \text{su} \\
 & d \alpha_i = -\epsilon_{abc} \alpha_b \wedge \alpha_c. \\
 & \alpha_i \in H^1 T^3 \cong \mathbb{R}^3. \\
 & T^3 \cong \mathbb{R}^3 / \mathbb{Z}^3, \\
 & \alpha_i = du_i, \quad i = 1, 2, 3. \\
 & \mathcal{A} \in H^3 M \cong \mathbb{R}, \quad \mathcal{B} \in H^4 M \cong \mathbb{R}^3. \\
 & \text{real} \quad m, n, k_1, k_2, k_3 \\
 & \mathcal{A} \in H^3 M \cong \mathbb{R} \quad m \\
 & \text{SU} \quad \rho \in \mathcal{A} \\
 & \rho = n_1 \alpha_1 + n_2 \alpha_2 + n_3 \alpha_3 + x_1 d_1 \alpha_1 + x_2 d_2 \alpha_2 + x_3 d_3 \alpha_3. \\
 & x_i t \quad t \quad m \quad n \quad t \quad \rho
 \end{aligned}$$

OV Z LO

$$\sigma \begin{matrix} k_1 & 1 & 2 & 3\alpha_1 & k_2 & 1 & 2 & 3\alpha_2 & k_3 & 1 & 2 & 3\alpha_3 \\ y_1 & 2\alpha_2 & 3\alpha_3 & & y_2 & 3\alpha_3 & 1\alpha_1 & & y_3 & 1\alpha_1 & 2\alpha_2 & \end{matrix}$$

$$\begin{matrix} 2\alpha_2 & 3\alpha_3 & d & 1\alpha_2\alpha_3 & d & 1\alpha_2\alpha_3, \\ & & & \rho & & \sigma \\ & & & k_{1,2,3} & & \\ \sigma & & & & & \\ & & & & & V^2 \sigma & y_1 y_2 y_3. \end{matrix}$$

$$\begin{matrix} k_1 & k_2 & k_3 & & SU & & \sigma \\ \sigma & y_1 d & 1\alpha_2\alpha_3 & y_2 d & 2\alpha_3\alpha_1 & y_3 d & 3\alpha_1\alpha_2. \\ & & \sigma & & & \omega^2/ & \omega \end{matrix}$$

$\omega \wedge \rho$

$$\omega \begin{matrix} \sqrt{\frac{y_2 y_3}{y_1}} & 1\alpha_1 & \sqrt{\frac{y_1 y_3}{y_2}} & 2\alpha_2 & \sqrt{\frac{y_1 y_2}{y_3}} & 3\alpha_3. \\ & & \omega & & & \\ & & \sigma & -\omega^2, & & \end{matrix}$$

$$\omega \wedge \rho \quad .$$

fl

$$x_i t \quad y_i t$$

$$\langle 2\alpha_2 \ 3\alpha_3, d \ 1\alpha_1 \rangle \int_{S^3 \times T^3} 2\alpha_2 \ 3\alpha_3 \ 1\alpha_1 \quad S^3 \quad T^3 / .$$

$$SU \quad \times \quad SU$$

$$dx_1 \wedge dy_1 \quad dx_2 \wedge dy_2 \quad dx_3 \wedge dy_3.$$

$$V \rho$$

$$V^2 \rho \quad -m^2 n^2 - m x_1 x_2 x_3.$$

$$V \rho \quad V \sigma$$

$$y_1 y_2 y_3 > , \quad m x_1 x_2 x_3 < -m^2 n^2.$$

$$H = \frac{V_\rho - V_\sigma}{\sqrt{-m^2 n^2 - m x_1 x_2 x_3} - \sqrt{y_1 y_2 y_3}},$$

$$i / j / k / i \quad , \quad , \quad , \quad \text{fl}$$

$$y_i = \frac{m x_j x_k}{\sqrt{-m^2 n^2 - m x_i x_j x_k}},$$

$$x_i = -\sqrt{\frac{y_j y_k}{y_i}}.$$

$$S^3 \times T^3 \quad G_2 \quad a, b \times$$

$$dt \wedge \left(\sqrt{\frac{y_2 y_3}{y_1}} \alpha_1 + \sqrt{\frac{y_1 y_3}{y_2}} \alpha_2 + \sqrt{\frac{y_1 y_2}{y_3}} \alpha_3 \right)$$

$$n \alpha_1 \alpha_2 \alpha_3 - m x_1 x_2 x_3 \quad x_1 d\alpha_1 + x_2 d\alpha_2 + x_3 d\alpha_3 .$$

$$X \quad a, b \times SU(n) \times T^3 \quad B \quad SU \quad a, b \times T^3$$

$$S^3 \quad T^3 \quad G_2$$

$$S^3 \geq |n|,$$

$$T^3 \geq |m|.$$

$$G_2$$

$$ds^2 = K_\Phi^{-1/9} \left(\sqrt{y_1 y_2 y_3} dt^2 + \sqrt{\frac{y_2 y_3}{y_1}} x_2 x_3 \alpha_1^2 + mn \alpha_1 - m x_1 \alpha_1^2 \right.$$

$$\left. + \sqrt{\frac{y_1 y_3}{y_2}} x_1 x_3 \alpha_2^2 + mn \alpha_2 - m x_2 \alpha_2^2 + \sqrt{\frac{y_1 y_2}{y_3}} x_1 x_2 \alpha_3^2 + mn \alpha_3 - m x_3 \alpha_3^2 \right)$$

$$K_\Phi = -m^3 y_1 y_2 y_3^{3/2} (x_1 x_2 x_3 - mn^2)^3 \quad y_1 y_2 y_3^{9/2},$$

$$K_\Phi^{-1/9} \quad H \quad dt^2$$

$$ds^2 = dt^2 + \frac{1}{y_1} x_2 x_3 \alpha_1^2 + mn \alpha_1 - m x_1 \alpha_1^2$$

$$+ \frac{1}{y_2} x_1 x_3 \alpha_2^2 + mn \alpha_2 - m x_2 \alpha_2^2 + \frac{1}{y_3} x_1 x_2 \alpha_3^2 + mn \alpha_3 - m x_3 \alpha_3^2 .$$

4.3. SU Symmetric Solution and Large Distance Asymptotics

$$\begin{aligned}
 & G_2 \\
 & m \quad n \\
 & x_1 \quad x_2 \quad x_3, \quad y_1 \quad y_2 \quad y_3. \\
 & SU \qquad \qquad \qquad SU
 \end{aligned}$$

f.

$$\begin{aligned}
 x &= -\sqrt{y}, \\
 y &= \frac{mx^2}{\sqrt{-m^2n^2 - mx^3}}. \\
 x t & \quad y t \\
 -\infty < x t & \leq -\left(\frac{mn^2}{m}\right)^{1/3}, \\
 & \leq y t < \infty.
 \end{aligned}$$

$H t$

$$-mx^3 \quad y^3 \quad -m^2n^2.$$

$$\frac{dy}{dt} = \frac{m}{y^{3/2}} \left(-y^3 - mn^2 \right)^{2/3},$$

$$t - t_0 = -\left(\frac{7}{m^5n^4}\right)^{1/3} y^{5/2} F\left(\left[-, -\right], \left[-\right], -\frac{y^3}{m^2n^2}\right).$$

G_2

$$ds^2 = dt^2 - \sum_{j=1}^3 x_j^2 \left(\frac{2}{y} mn_j \alpha_j - mx \alpha_j^2 \right),$$

$$SU \times SU \times U^3.$$

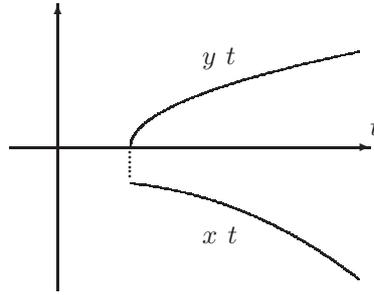


Figure 2.

$x(t)$ $y(t)$

$x(t)$ $y(t)$

Solution with n and Large Distance Asymptotics

n

$$x = -\frac{m^{1/3}}{2} (t - t_0)^2, \quad y = \frac{m^{2/3}}{2} (t - t_0)^2.$$

t_0

$$ds^2 = dt^2 - t^2 \left(\frac{2}{1} + \frac{2}{2} + \frac{2}{3} \right) m^{2/3} \alpha_1^2 \alpha_2^2 \alpha_3^2.$$

/ i

fl

$$\mathbb{R}^4 \times T^3.$$

m

T^3

$$T^3|_{t=\infty} = m,$$

SU

m

fl

m

n

$y(t)$

OV Z LO

$$\begin{aligned} x_1 &= -\sqrt{y_3}, & y_1 &= \frac{mx_1x_3}{y_1\sqrt{y_3}}, \\ x_3 &= -\frac{y_1}{\sqrt{y_3}}, & y_3 &= \frac{mx_1^2}{y_1\sqrt{y_3}}. \end{aligned}$$

n

$$t = t_0 - x_1 \left(-\frac{B^2}{\beta m} \right)^{1/4} F \left(\left[- , - \right], \left[- \right], -Bx_1^2/\beta \right),$$

$$B \leq \frac{\beta}{x_3 y_3 - x_1 y_1} B \quad SU$$

$t \rightarrow$

$$x_i = y_i t \quad S^3$$

$$S^3 \rightarrow S^2, \quad t \rightarrow .$$

$$x_1 t = y_1 t \quad t \quad x_3 t = y_3 t$$

$\beta <$

$$y_3 \approx \frac{\beta}{x_3},$$

$$y_1 \approx \sqrt{-\frac{m}{\beta}} x_1 x_3$$

$x_1 \quad x_3$

$$\begin{aligned} x_1 &= -A \gamma t, & y_1 &= \frac{\sqrt{-m\beta}}{\gamma^2 A^2} \gamma t, \\ x_3 &= \frac{\beta}{\gamma^2 A^2} \gamma t, & y_3 &= \gamma^2 A^2 \gamma t, \end{aligned}$$

$$\begin{aligned} & \gamma^2 \sqrt{\frac{m}{|\beta|}} \\ & |\beta| \gg |x_1 y_1| \\ & \gamma t \ll t \end{aligned}$$

$$\begin{aligned} x_1 & \approx -A\gamma t, & y_1 & \approx \frac{\sqrt{-m\beta}t}{\gamma A}, \\ x_3 & \approx \frac{\beta}{\gamma^2 A^2}, & y_3 & \approx \gamma^2 A^2. \end{aligned}$$

$$\begin{aligned} ds^2 & \approx dt^2 \sqrt{\frac{|\beta|}{m}} \left(\frac{1}{2} \alpha_1^2 + \frac{1}{2} \alpha_2^2 + \frac{mA^2}{|\beta|} \alpha_1^2 \alpha_2^2 + t^2 \frac{2}{3} \frac{\beta^2}{A^4} \alpha_3^2 \right) \\ & \quad \gamma t \gg G_2 \quad \mathbb{R}^4/\mathbb{Z}_2 \times T^3 \\ X & \sim S^2 \times \mathbb{R}^2 \times T^3, \quad \mathbb{R}^4/\mathbb{Z}_2 \times T^3 \\ & \quad \sqrt{|\beta|} \end{aligned}$$

4.5. General Solution

$$\begin{aligned} H & \sqrt{-m^2 n^2 - mx_1 x_2 x_3} - \sqrt{y_1 y_2 y_3} \quad V \rho - V \sigma, \\ H & , \quad V \rho \quad V \sigma, \\ \tilde{H} & V \rho^2 - V \sigma^2 - m^2 n^2 - mx_1 x_2 x_3 - y_1 y_2 y_3. \\ dV \rho^2 - V \sigma^2 & \quad V \rho \quad dV \rho - dV \sigma \end{aligned}$$

$$\begin{aligned} \frac{dt}{\tilde{t}} & \quad V \rho \quad V \sigma \quad \sqrt{y_1 y_2 y_3}. \\ \vec{x}, \vec{y} & \mapsto M \cdot \vec{x}, M^{-1} \cdot \vec{y}, \quad M \end{aligned}$$

$$\begin{aligned} & x_1 y_1 - x_3 y_3 \quad x_2 y_2 - x_3 y_3. \\ i / j / k / i. & \\ & x_i - y_j y_k \quad y_i \quad m x_j x_k \end{aligned}$$

OV Z LO

$\frac{d}{dt}$.

$$z_i = x_i y_i.$$

$$\frac{d}{dt} z_i - z_j = \dots,$$

$$\begin{aligned} z &\equiv z_1 \\ z_2 &= z \alpha \\ z_3 &= z \beta. \end{aligned}$$

$$X \equiv x_1 x_2 x_3 \quad Y \equiv y_1 y_2 y_3.$$

$$m^2 n^2 = mX - Y.$$

$$z = x_1 y_1 = x_2 y_2 = x_3 y_3 = Y/mX = m^2 n^2 / Y.$$

$$Y = m x_2 x_3 y_2 y_3 = x_1 x_3 y_1 y_3 = x_1 x_2 y_1 y_2,$$

$$\ddot{z} = m z^2 = z \alpha \beta = \alpha \beta.$$

$$Y = z m^2 n^2 / m,$$

$$x_i = y_i = x_1 = -y_2 y_3 = -Y/y_1 = -Y/z x_1,$$

$$x_1 = -A_1 = -\int Y/z \, d\tilde{t},$$

$$A_1 = y_1 = z/x_1. \quad x_2, y_2, x_3, y_3$$

$t,$

$$\frac{dt}{d\tilde{t}} = \sqrt{Y}.$$

Equation for z

$$u = mz - m\alpha - \beta$$

$$\ddot{u} = u^2 - D$$

$$D = m^2(\alpha\beta - \alpha^2 - \beta^2) - m^2(\alpha^2 - \alpha\beta - \beta^2)$$

$$u = f(u)$$

$$\ddot{u} = f'u - ff'$$

$$f^2 = \frac{2}{3}u^3 - Du - E$$

$$f^2 = u^2 - D$$

$$\tilde{t} = C \pm \int \frac{du}{\sqrt{\frac{2}{3}u^3 - Du - E}}$$

$$y^2 = -x^3 - Dx - E$$

$$u = \tilde{t}$$

$$v = u$$

$$g_2 \equiv -D$$

$$v$$

$$\ddot{v} = v^2 - g_2$$

$$\tilde{t} = C \pm \int \frac{dv}{\sqrt{v^3 - g_2v - g_3}}$$

$$C = g_3$$

$$v = p_\tau \tilde{t} - C$$

$$\tau$$

$$g_2$$

$$g_3$$

$$-\tilde{t} - C$$

$$p$$

$$p'^2 = p^3 - g_2p - g_3$$

OV Z LO

$$p'p'' = p^2p' - g_2p',$$

$$p'' = p^2 - g_2 / .$$

$g_2 \quad g_3$

$$v = \frac{g_2}{\tilde{t}^2} C ,$$

$g_2 /$

$$v = \frac{g_2}{\tilde{t}^2} C - / .$$

Asymptotics and Behavior Near the Poles

p

$$\mathbb{R}^4 \times T^3$$

$$v \approx \tilde{t}^{-2}.$$

$$v = \frac{g_2}{\tilde{t}^2}$$

$$z \approx -\frac{u}{m} \approx -\frac{v}{m} \approx -\frac{v}{m\tilde{t}^2}.$$

$$x_i \tilde{t} \quad y_i \tilde{t}$$

$$x_i \sim y_i \sim \frac{m}{\tilde{t}},$$

$$Y \approx \frac{m}{m\tilde{t}^3}.$$

$$t = \tilde{t}$$

t

$$x_1 = -A_1 t^2, \quad y_1 = \frac{m}{A_1} t^2,$$

$$x_2 = -A_2 t^2, \quad y_2 = \frac{m}{A_2} t^2,$$

$$x_3 = -A_3 t^2, \quad y_3 = \frac{m}{A_3} t^2,$$

OV Z LO

$x_i <$

A_i

$$A_1 A_2 A_3 = m.$$

$$A_3 = \frac{A_i A_3}{A_1 A_2}.$$

x_i

$n \rightarrow$

$$y_i = \sqrt{-\frac{m x_j x_k}{x_i}}, \quad x_i = -\sqrt{\frac{y_j y_k}{y_i}}.$$

SU

$x_i \quad y_i$

t^2

G_2

$$ds^2 = dt^2 - t^2 \left(\frac{dx_1^2}{x_1^2} + \frac{dx_2^2}{x_2^2} + \frac{dx_3^2}{x_3^2} + A_1^2 \alpha_1^2 + A_2^2 \alpha_2^2 + A_3^2 \alpha_3^2 \right).$$

\mathbb{R}^4

/

$$\frac{x_1 x_2}{t^2 y_3} = A_1 A_2 \left(\frac{m}{A_3} \right)^{-1} = \frac{A_1 A_2 A_3}{m} = -.$$

A
fl

$\mathbb{R}^4 \times T^3.$

SU

$$T^3 = \int_{T^3} \sqrt{g} = \int_{T^3} \sqrt{3/2} A_1 A_2 A_3 = m,$$

$\mathbb{R}^4 \times T^3$

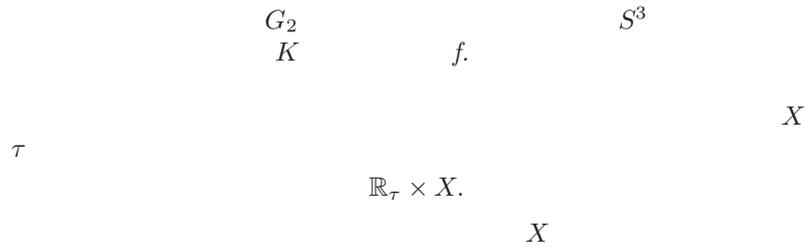
$m \quad n$

t

OV Z LO



5. Abelian BPS Monopoles from Torus Fibrations



¹⁴Kaluza-Klein reduction of certain G_2 holonomy metrics to non-abelian monopoles has been discussed recently in [13].

G_2	X	
SU		
U		

$SU \quad U$

$\tau \quad t \quad r \quad t$

$\mathbb{R}_\tau \times X$

$$ds^2 = -d\tau^2 + dr^2 + \frac{x_2 x_3}{y_1} \alpha_1^2 + mn_1 \alpha_1 - mx_1 \alpha_1^2$$

$$- \frac{x_1 x_3}{y_2} \alpha_2^2 + mn_2 \alpha_2 - mx_2 \alpha_2^2 - \frac{x_1 x_2}{y_3} \alpha_3^2 + mn_3 \alpha_3 - mx_3 \alpha_3^2,$$

$x_i \quad y_i \quad r$

$x_i \equiv x_i r,$
 $y_i \equiv y_i r.$

$G_2 \quad X \quad SU \times T^3$
 $U \quad U \subset SU$

$\phi \quad f.$
 T^3

$\alpha_4 \equiv d\phi.$

$\alpha_i \quad i \quad , \quad , \quad ,$

$T^4 \quad T^3 \times U$
 T^4

$$ds^2 = ds_{1,3}^2 + h_{ij} \alpha_i A_i \alpha_j A_j,$$

$ds_{1,3}^2$
 $A_i \quad i \quad U$
 $h_{ij} \quad i \quad U \quad - \quad j \quad U$

4D Theory:

$$h \begin{pmatrix} -\frac{mx_1}{y_1} & & & \frac{mn}{2y_1} & \psi & \theta \\ & -\frac{mx_2}{y_2} & & \frac{mn}{2y_2} & \psi & \theta \\ & & -\frac{mx_3}{y_3} & \frac{mn}{2y_3} & \theta & \\ \frac{mn}{2y_1} & \psi & \theta & \frac{mn}{2y_2} & \psi & \theta \\ & & & \frac{mn}{2y_3} & \theta & \\ & & & & & h_{44} \end{pmatrix},$$

$$h_{44} = \frac{x_1 x_2}{y_3} - 2\theta - \frac{x_3}{y_1 y_2} x_1 y_1 - 2\psi - x_2 y_2 - 2\psi - 2\theta.$$

$$A_i = \sum_{k=1}^4 h_{ik}^{-1} \tilde{A}_k,$$

$$\begin{aligned} \tilde{A}_1 &= \frac{mn}{y_1} \psi d\theta, \\ \tilde{A}_2 &= -\frac{mn}{y_2} \psi d\theta, \\ \tilde{A}_3 &= \frac{mn}{y_3} d\psi, \\ \tilde{A}_4 &= \frac{x_1 x_2}{y_3} \theta d\psi - \left(\frac{x_2 x_3}{y_1} - \frac{x_1 x_3}{y_2} \right) \psi - \psi - \theta d\theta. \end{aligned}$$

$$ds_{1,3}^2 = -d\tau^2 - dr^2 - \frac{x_1 x_2}{y_3} d\psi^2 - \left(\frac{x_2 x_3}{y_1} - 2\psi - \frac{x_1 x_3}{y_2} - 2\psi \right) d\theta^2 - \sum_{i,j=1}^4 A_i h_{ij}^{-1} A_j.$$

5.1. Spherically Symmetric Monopoles

 SU

$$x_1, x_2, x_3, y_1, y_2, y_3.$$

$$A_1 = -\left(\frac{n}{x} \right) \psi d\theta - \psi - \psi - \theta d\psi,$$

OV Z LO

$$\begin{aligned}
 A_2 & \left(\frac{n}{x}\right) \psi d\theta \quad \psi \quad \theta \quad \theta d\psi, \\
 A_3 & -\left(\frac{n}{x}\right)^2 \theta d\psi, \\
 A_4 & \theta d\psi.
 \end{aligned}$$

$A_1 \ A_2 \ A_3 \ SU \ n \ A_4$

$$ds_E^2 = \left(-\frac{mx}{y^{1/2}}\right) -d\tau^2 - dr^2 - y^{3/2} d\Omega_2^2.$$

$\mathbb{R}^4 \times T^3 \ n \ , \ G_2 \ x \ r \ y \ r \ X \ fl$

$$x = -\frac{m^{1/3}}{r^2}, \quad y = \frac{m^{2/3}}{r^2}.$$

$$h = m^{2/3}, m^{2/3}, m^{2/3}, \varphi, \quad \varphi = -r^2,$$

$$ds_E^2 = -mr -d\tau^2 - dr^2 - \frac{m}{r^3} d\Omega_2^2.$$

A_3

$A_1 \ A_2$

$U \ A_4$

$$A_4 \quad \theta d\psi.$$

A_4

φ

$$S = -\frac{1}{\pi \kappa_5} \int d^5x \sqrt{-g_5} R_5.$$

φ

\mathbb{R}^3

ϕ

f .

$x_i x_j / y_k$

r

OV Z LO

${}^4 \times T^3$

\mathbb{R}^4

fl $\mathbb{R}^4 \times T^3$

SU

T^3

5.2. Axially Symmetric Monopoles

U x_1 x_2 y_1 y_2 U SU

$$A_4 \frac{z}{\beta} \frac{\theta}{z^2 \theta} d\psi,$$

ψ z $x_1 y_1$ A_4

$$ds_E^2 = \sqrt{h} \left[-d\tau^2 - dr^2 - \frac{y_1^2 y_3}{mz} \left(d\theta^2 - \frac{z}{\beta} \frac{\theta}{z^2 \theta} d\psi^2 \right) \right]$$

$$- \frac{m \sqrt{z^2 - \beta z \theta}}{y_1 \sqrt{y_3}} \left(-d\tau^2 - dr^2 - \frac{y_1^2 y_3}{mz} d\theta^2 \right) - \frac{y_1 \sqrt{y_3} z \theta}{\sqrt{z^2 - \beta z \theta}} d\psi^2.$$

$|z| \gg |\beta|$

θ

$SO \sim SU$

U

$\beta \rightarrow$

$$ds_E^2 \approx \frac{m}{\sqrt{\theta}} \left(-d\tau^2 - dr^2 - \frac{|\beta|}{m} d\theta^2 + \sqrt{\frac{|\beta|}{m}} r^2 d\psi^2 \right).$$

Acknowledgments

References

[] “E ” Winter School on Mirror Symmetry,
Vector Bundles and Lagrangian Submanifolds, V

[] “O E H ” N
B524 8

[] E “ F 2 H ”
 / 5

[4] / 5 V “ fl N ”

[5] E “ 2 ”
 /

[] V V “ F ” N
B463 4

[] “ N

[8] R “O / 4 E
 H ” *Duke Math. J.* **58** 8 8
 N N “E $S^3 \mathbb{R}^3 \mathbb{R}^4$ ”
Commun.Math.Phys **127** 5 55

[] R “ 2 ” / 5
 [] V “ N

[] “ N **B337**
 H “L L ”

[] / 8 8
 [] “ ” N **B574**

[] H “ ” / 5
 [4] H “ H ” N **B449** 5
 5 5 55 ; — **B458** 45
 [5] N H “ 4 **25** 5 5 5; L ” N
 [] N H “ ” / 54
 [] N H “ ” /
 [8] ” H ” O
 [] V “E R N = H ”
 N **B450** 5 8
 [] R “ E ”
Advanced Studies in Pure Mathematics 18-II: Kähler Metric and Moduli Spaces, O
 5
 [] “ R ”
 [] H L / 8 N L “ 2 *Spin* ”
 / 45
 [] “O R
 ” N **B593** ; / 8
 [4] R L “ ” **6** 8
 5 4
 [5] N H E ” N
 H ” N **B279** 8
 [] H O V “ ” R L **77** 8
 [] H O O Z “ ” N
B477 4 4
 [8] “ =
 ” *Phys.Lett.* **B357** 5
 [] *J. Phys.* **16** 8 5 5; *Nucl. Phys.* **B226** 8 ;
 R *Phys. Rev. Lett.* **51** 8 8
 [] “H ” N **B153**
 [] L V “ E H ”
 / 4 5
 [] E Z “ ” N
B479 4 5
 [] V “RR F ”
 L **B474**
 [4] “ H ”
 [5] V E “ N **6** 8 5 8 ” N **B431** 4

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DEPARTMENT OF PHYSICS, HARVARD UNIVERSITY, CAMBRIDGE, MA 02136, USA
E-mail address: gukov@democritus.harvard.edu

DEPARTMENT OF MATHEMATICS, HARVARD UNIVERSITY, CAMBRIDGE, MA 02138, USA
E-mail address: yau@math.harvard.edu

DEPARTMENT OF MATHEMATICS, NORTHWESTERN UNIVERSITY, 2033 SHERIDAN ROAD, EVANSTON, IL
60208-2730, USA
E-mail address: zaslow@math.northwestern.edu