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Research Article

## **On operator systems generated by reducible projective unitary representations of compact groups**

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**Abstract:** We study reducible projective unitary representations  $(U_g)_{g \in G}$  of a compact group *G* in separable Hilbert spaces *H*. It is shown that there exist the projections *Q* and *P* for which  $V = span(U_g Q U_g^*$ ,  $g \in G$  is the operator system and  $PVP = \{CP\}$ . As an example, a bipartite Hilbert space  $H = 5 \otimes 5$  is considered. In this case, the action of  $(U_q)_{q \in G}$  has the property of transforming separable vectors to entangled.

**Key words:** Operator systems, covariant resolutions of identity, reducible unitary representations of compact groups, quantum anticliques

#### **1. Introduction**

A subspace *V* consisting of bounded linear operators in a separable Hilbert space *H* is said to be an operator system [\[5](#page-5-0)] if it is self-adjoint  $(V \in V)$  implies  $V^* \in V$  and the identity operator  $I \in V$ . Recently, operator systems have attracted the interest of researchers in the context of both functional analysis and quantum information theory  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$  $[1-4, 6, 7, 9, 10, 12]$ . It should be noted that operator systems are often called noncommutative operator graphs.

The Kraus representation of a quantum channel generates the operator system [[7\]](#page-5-2). The possibility of transmitting quantum information via a channel with zero error is completely determined by the properties of the operator system corresponding to this channel  $[3, 4, 9, 10]$  $[3, 4, 9, 10]$  $[3, 4, 9, 10]$  $[3, 4, 9, 10]$  $[3, 4, 9, 10]$  $[3, 4, 9, 10]$  $[3, 4, 9, 10]$ . Moreover, it is hoped that the proximity of the two quantum channels [[11\]](#page-5-6) can be estimated using the corresponding operator systems.

Let *G* be a compact group with the Haar measure  $\mu$ ,  $\mu(G) = 1$ , and **B** is the sigma-algebra generated by compact subsets of *G*, then the set of positive operators  $\{M(B), B \in \mathfrak{B}\}\$  in a Hilbert space *H* is said to be a resolution of identity if [\[8](#page-5-7)]

$$
M(\emptyset) = 0, \ M(G) = I,
$$

$$
M(\cup_j B_j) = \sum_j M(B_j), \ B_k \cap B_l = \emptyset \text{ for } k \neq l, \ B_j \in \mathfrak{B},\tag{1.1}
$$

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and a convergence in  $(1.1)$  $(1.1)$  $(1.1)$  is understood in the sense of strong operator topology. Let  $(U_g)_{g \in G}$  be a projective unitary representation of *G* in *H*. Then  $\{M(B), B \in \mathfrak{B}\}\$ is said to be covariant with respect to the action of  $(U_q)_{q∈G}$  if

<span id="page-2-0"></span>
$$
U_g M(B) U_g^* = M(gB).
$$

In [\[2](#page-4-3)], the study of operator systems generated by covariant resolutions of identity in the sense of

$$
\mathcal{V} = \overline{span\{M(B), B \in \mathfrak{B}\}}\tag{1.2}
$$

was initiated. It is known [\[8](#page-5-7)] that in finite dimensional spaces *H*, any covariant resolution of identity  ${M(B), B \in \mathfrak{B}}$  has the form

<span id="page-2-1"></span>
$$
M(B) = \int\limits_B U_g M_0 U_g^* d\mu(g),
$$

where  $M_0$  is some positive operator in  $H$ . In this case,  $(1.2)$  $(1.2)$  $(1.2)$  can be replaced with

$$
\mathcal{V} = \overline{span(U_g M_0 U_g^*, \ g \in G)}.\tag{1.3}
$$

A particularly interesting case is the bipartite Hilbert space  $H = \mathfrak{H} \otimes \mathfrak{H}$ . Then, a vector  $v \in H$  is said to be separable if it can be represented in the form  $v = v_1 \otimes v_2$ . In the opposite case, *v* is known as entangled. Given a unit vector  $v \in H$  denote  $|v\rangle \langle v|$  an orthogonal projection to the subspace  $\{Cv\}$ . We call a projection *Q* to be separable if it can be represented as a sum of orthogonal projections  $|v\rangle\langle v|$  for which  $v$  are separable vectors. In [[3,](#page-4-2) [4\]](#page-4-1) operator systems generated by unitary representations of the circle group as well as the Heisenberg–Weyl group *G* were studied in detail. Operator systems were constructed having the form  $(1.2)$  $(1.2)$  $(1.2)$  and  $(1.3)$  for which  $M_0 = Q$  are separable projections. Moreover it was proved that for such *V* there are projections *P* satisfying the property

$$
PVP = \{ \mathbb{C}P \}. \tag{1.4}
$$

<span id="page-2-2"></span>The projections *P, rankP*  $\geq$  2, satisfying [\(1.4](#page-2-2)) are known as quantum anticliques for  $V$  [[12\]](#page-5-5). In the present paper, how these ideas work in general case is shown.

#### **2. Reducible projective unitary representation of groups in an arbitrary separable Hilbert space**

In this section, we do not take into account a tensor product structure of Hilbert space *H* . Consider a reducible continuous projective unitary representation  $G \ni g \to U_g$  in *H*. Then, there are countably many resolutions of *H* into the orthogonal sum

<span id="page-2-3"></span>
$$
H = \bigoplus_j H_j,\tag{2.1}
$$

such that the restrictions

<span id="page-2-4"></span>
$$
U_g^{(j)} = U_g|_{H_j} \tag{2.2}
$$

determine cyclic representations of *G*. It means that for any *j*  $U_q H_j \subset H_j$ ,  $g \in G$ , and there exist a unit vector  $v_j \in H_j$  such that the closer  $\overline{span(U_gv_j, g \in G)} = H_j$ . Since all irreducible representations of compact group are finite dimensional, the following statement holds true.

**Proposition 2.1** *Suppose that at least one of the following conditions is satisfied*

- *• all representations [\(2.1\)](#page-2-3) and ([2.2\)](#page-2-4) are irreducible;*
- *• dimH<sup>j</sup> <* +*∞ and cyclic vectors v<sup>j</sup> has the property*

$$
\int_G U_g \left| v_j \right\rangle \left\langle v_j \right| U_g^* d\mu(g) = \frac{1}{dim H_j} Q_j,
$$

*where*  $Q_j$  *is an orthogonal projection on*  $H_j$ .

∫

*G*

*Then, the positive operator*

<span id="page-3-2"></span>
$$
Q = \sum_{j} dim H_{j} |v_{j}\rangle\langle v_{j}|
$$
\n(2.3)

*generates an operator system by the formula*

<span id="page-3-0"></span>
$$
\mathcal{V} = \overline{span(U_g Q U_g^*, \ g \in G)}.\tag{2.4}
$$

**Proof** It immediately follows that *V* is self-adjoint. Hence, it remains to prove that  $I \in \mathcal{V}$ . Since  $dim H_j < +\infty$ , we get

$$
U_g^{(j)} |v_j\rangle \langle v_j | U_g^{(j)*} d\mu(g) = \frac{1}{dim H_j} Q_j.
$$
  

$$
\int_G U_g Q U_g^* d\mu(g) = I.
$$

Thus,

Now suppose that the spectral decompositions of  $U_g^{(j)}$  contain the same projection  $|h_j\rangle \langle h_j|$  for all  $g \in G$ . Recall that two vectors  $v$  and  $h$  in a finite dimensional space  $K$  are said to be unbiased if

$$
|\langle v, h \rangle|^2 = \frac{1}{dim K}.\tag{2.5}
$$

**Proposition 2.2** *Suppose that vectors*  $h^j$  *and*  $v_j$  *are unbiased for all j*. Then, the projection

$$
P=\sum_j\ket{h^j}\bra{h^j}
$$

*is a quantum anticlique for ([2.4\)](#page-3-0).*

**Proof** It follows from  $(2.5)$  that

$$
\left|h^{j}\right\rangle\left\langle h^{j}\right|\left|v_{j}\right\rangle\left\langle v_{j}\right|\left|h^{j}\right\rangle\left\langle h^{j}\right|=\frac{1}{dimH_{j}}\left|h^{j}\right\rangle\left\langle h^{j}\right|.
$$

Taking into account [\(2.3\)](#page-3-2), we get

$$
PQP = P
$$

because  $U_g P = P U_g$  are commuting for all  $g \in G$ .

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<span id="page-3-1"></span> $\Box$ 

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#### **3. Operator systems in a bipartite Hilbert space**

Here we shall give an explicit example showing how the techniques of the previous section works in a bipartite finite dimensional Hilbert space  $H = \mathfrak{H} \otimes \mathfrak{H}$ . Denote by  $(|jk\rangle)$  the basis in *H* consisting of separable vectors,  $1 \leq j, k \leq dim\mathfrak{H} < +\infty$ . Together with  $(|jk\rangle)$ , we consider the basis in *H* consisting of entangled vectors

$$
h_k^j = \frac{1}{\sqrt{d}} \sum_{s=1}^d e^{i\frac{2\pi ks}{d}} \left| s\ s + j\ mod\ dim \mathfrak{H} \right\rangle.
$$

Denote by  $H_j$  the subspaces spanned by vectors  $h_k^j$ ,  $1 \leq k \leq dim \mathfrak{H}$ .

In the following statement we claim that the conditions of Proposition 1 are satisfied.

**Theorem 3.1** Fix  $1 \leq m_0, n_0 \leq dim\mathfrak{H}$ . Suppose that  $G \ni g \to U_g$  can be resolved in a sum of cyclic *representations*

$$
U_g=\oplus_j U_g|_{H_j}
$$

with the cyclic vectors  $v_j^{m_0} = |m_0 m_0 + j \mod{dim \mathfrak{H}}\rangle$  and the projections  $|h_{n_0}^j\rangle\langle h_{n_0}^j|$  are contained in the *spectral decompositions of*  $U_g$  *for all*  $g \in G$ *. Then, the projection* 

<span id="page-4-4"></span>
$$
P=\sum_j\ket{h_{n_0}^j}\bra{h_{n_0}^j}
$$

*is a quantum anticliques for the operator system*

$$
\mathcal{V} = \overline{span(U_g Q U_g^*, \ g \in G)},\tag{3.1}
$$

*where*

$$
Q = \sum_j \left| v_j^{m_0} \right\rangle \left\langle v_j^{m_0} \right|.
$$

**Remark 3.2** *Since the projections*  $U_g Q U_g^*$  *have infinite ranks in general, we need to take a closer look at [\(3.1\)](#page-4-4) in the sense of strong operator topology to guarantee the inclusion*  $I \in \mathcal{V}$ *.* 

**Proof** Proposition 1 implies that  $(3.1)$  $(3.1)$  is an operator system. Now it is enough to check that  $h_{n_0}^j$  and  $v_j^{m_0}$ are unbiased and the result follows from Proposition 2.  $\Box$ 

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